# Exponential decay estimates for fundamental solutions of Schrödinger-type operators

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### The electric Schrödinger operator

By an electric Schrödinger operator we mean a second-order linear operator of the form

 $L := -\mathsf{div}A\nabla + V,$ 

where V is a scalar, real-valued, positive a.e., locally integrable function, and A is a matrix of bounded, measurable complex coefficients satisfying the uniform ellipticity condition

$$\lambda |\xi|^2 \le \Re e \langle A(x)\xi,\xi\rangle \equiv \Re e \sum_{i,j=1}^n A_{ij}(x)\xi_j \bar{\xi_i} \text{ and } \|A\|_{L^{\infty}(\mathbb{R}^n)} \le \Lambda,$$
(1)

#### for some $\lambda > 0$ , $\Lambda < \infty$ and for all $\xi \in \mathbb{C}^n$ , $x \in \mathbb{R}^n$ .

• The exponential decay of solutions to the Schrödinger operator in the presence of a positive potential is an important property underpinning foundation of quantum physics.

• However, establishing a precise rate of decay for complicated potentials is a challenging open problem to this date. (Landis conjecture)

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## Exponential Decay of electric Schrödinger operators

## • First results expressing upper estimates on the solutions in terms of a certain distance associated to V go back to Agmon [1], but not sharp.

• For eigenfunctions, the decay is governed by the uncertainty principle - see ADFJM '*Localization of eigenfunctions via an effective potential*"[2]

• In [9], Shen proved that if  $V \in RH_{\frac{n}{2}}$ , then the fundamental solution  $\Gamma$  to the classical Schrödinger operator  $-\Delta + V$  satisfies the bounds

$$\frac{c_1 e^{-\varepsilon_1 d(x,y,V)}}{|x-y|^{n-2}} \le \Gamma(x,y) \le \frac{c_2 e^{-\varepsilon_2 d(x,y,V)}}{|x-y|^{n-2}},$$
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• A natural question is whether the sharp exponential decay found by Shen for the fundamental solution to  $-\Delta + V$  can be extended to the non self-adjoint setting  $-\text{div }A\nabla + V$ .

• Moreover, we also wondered whether we can obtain exponential decay results for the fundamental solution to the *magnetic Schrödinger* operator.

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We consider the operator formally given by

$$L = -(\nabla - i\mathbf{a})^T A(\nabla - i\mathbf{a}) + V,$$
(3)

where  $\mathbf{a} = (a_1, \dots, a_n)$  is a vector of real-valued  $L^2_{loc}(\mathbb{R}^n)$  functions, A and V as before. Denote

$$D_{\mathbf{a}} = \nabla - i\mathbf{a},$$

and the magnetic field by  $\mathbf{B}$ , so that

$$\mathbf{B} = \operatorname{curl} \mathbf{a}. \tag{4}$$

## Properties of the magnetic Schrödinger operator

• The magnetic Schrödinger operator exhibits a property called *gauge invariance*: quantitative assumptions should be put on  $\mathbf{B}$  rather than  $\mathbf{a}$ .

• The diamagnetic inequality

$$\left|\nabla|u|(x)\right| \le \left|D_{\mathbf{a}}u(x)\right|.\tag{5}$$

• When  $A \equiv I$  so that  $L_M := L = (\nabla - i\mathbf{a})^2 + V$ , the operator  $L_M$  is dominated by the Schrödinger operator  $L_E := -\Delta + V$  in the following sense: for each  $\varepsilon > 0$ ,

 $|(L_M + \varepsilon)^{-1}f| \le (-\Delta + \varepsilon)^{-1}|f|,$  for each  $f \in H = L^2(\mathbb{R}^n).$  (6)

The above is known as the *Kato-Simon inequality*.

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The above is known as the *Kato-Simon inequality*.

• Some authors (Ben Ali [3], Kurata and Sugano [7]) showed properties of the fundamental solution to  $L_M$  under ad-hoc smoothness assumptions on the magnetic potential **a**.

• On the other hand, the natural setting to make sense of  $L_M$  in the weak sense requires only that  $\mathbf{a} \in L^2_{loc}(\mathbb{R}^n), V \in L^1_{loc}(\mathbb{R}^n)$ .

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## Some more history

• Before Shen obtained the sharp exponential decay result for the fundamental solution to  $-\Delta + V$ , he first obtained a polynomial decay result in [11].

• He later also obtained polynomial decay results for the magnetic Schrödinger operator  $L_M$  in [12].

• In [6], Kurata obtained exponential decay results for  $-\text{div }A\nabla + V$ and  $L_M$  by integrating certain heat kernel estimates. He obtained the bound

$$|\Gamma(x,y)| \leq \frac{Ce^{-\varepsilon(1+m(x,V+|\mathbf{B}|)|x-y|)^{\frac{2}{2k_0+3}}}}{|x-y|^{n-2}} \quad \text{ for a.e. } x,y \in \mathbb{R}^n,$$

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To state our results, we need just a few more definitions.

We say that  $w \in L^p_{loc}(\mathbb{R}^n)$ , with w > 0 a.e., belongs to the Reverse Hölder class  $RH_p = RH_p(\mathbb{R}^n)$  if there exists a constant C so that for any ball  $B \subset \mathbb{R}^n$ ,

$$\left(f_B w^p\right)^{1/p} \le C f_B w. \tag{7}$$

For a function  $w \in RH_p, p \geq \frac{n}{2},$  define the maximal function m(x,w) by

$$\frac{1}{m(x,w)} := \sup_{r>0} \Big\{ r : \frac{1}{r^{n-2}} \int_{B(x,r)} w \le 1 \Big\},\tag{8}$$

and the distance function

$$d(x, y, w) = \inf_{\gamma} \int_{0}^{1} m(\gamma(t), w) |\gamma'(t)| dt,$$
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where  $\gamma:[0,1]
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The function m measures the sum of the contributions of the kinetic energy  $\Re eAD_{\mathbf{a}}f\overline{D_{\mathbf{a}}f}$  and potential energy  $V|f|^2$ , and is intimately related to the uncertainty principle through the following estimate which is often known as the *Fefferman-Phong* inequality:

Suppose that  $\mathbf{a} \in L^2_{loc}(\mathbb{R}^n)^n$ , and moreover assume (12) (next slide). Then, for all  $u \in C^1_c(\mathbb{R}^n)$ ,

$$\int_{\mathbb{R}^n} m^2(x, V + |\mathbf{B}|) |u|^2 \, dx \le C \int_{\mathbb{R}^n} (|D_{\mathbf{a}}u|^2 + V|u|^2) \, dx, \qquad (10)$$

where C depends on the constants c, c' from (12) and on  $||V + |\mathbf{B}||_{RH_{\frac{n}{2}}}$ .

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## $L^2$ exponential decay

#### Theorem 1 (Mayboroda, P. 2018)

For any operator L given by (3) and for any  $f \in L^2(\mathbb{R}^n)$  with compact support, there exist constants  $\tilde{d}, \varepsilon, C > 0$  such that

$$\int_{\left\{x \in \mathbb{R}^{n} | d(x, \operatorname{supp} f, V + |\mathbf{B}|) \geq \tilde{d}\right\}} m\left(\cdot, V + |\mathbf{B}|\right)^{2} |u|^{2} e^{2\varepsilon d\left(\cdot, \operatorname{supp} f, V + |\mathbf{B}|\right)} \\
\leq C \int_{\mathbb{R}^{n}} |f|^{2} \frac{1}{m(x, V + |\mathbf{B}|)^{2}}, \quad (11)$$

where  $u := L^{-1}f$  (in a weak sense), provided that A is an elliptic matrix with complex bounded measurable coefficients, and

i) either  $\mathbf{a} = 0$  and  $V \in RH_{n/2}$ ,

ii) or, more generally,  $\mathbf{a}\in L^2_{loc}(\mathbb{R}^n)$ , V>0 a.e. on  $\mathbb{R}^n$ , and

$$\begin{cases} V + |\mathbf{B}| \in RH_{n/2}, \\ 0 \le V \le c \, m(\cdot, V + |\mathbf{B}|)^2, \\ |\nabla \mathbf{B}| \le c' \, m(\cdot, V + |\mathbf{B}|)^3. \end{cases}$$
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## $L^2$ exponential decay for the resolvent

An analogous estimate holds for the resolvent operator  $(I + t^2 L)^{-1}$ , t > 0:

$$\int m\left(\cdot,\mathscr{B}_t\right)^2 \left| (I+t^2L)^{-1}f \right|^2 e^{2\varepsilon d\left(\cdot,\operatorname{supp} f,\mathscr{B}_t\right)} \left\{ x \in \mathbb{R}^n | d(x,\operatorname{supp} f,\mathscr{B}_t) \ge \tilde{d} \right\}$$

$$\leq C \int_{\mathbb{R}^n} |f|^2 m\Big(\cdot, \mathscr{B}_t\Big)^2.$$

where  $\mathscr{B} := V + |\mathbf{B}| + \frac{1}{t^2}$ .

• In other words,  $L^{-1}f$  decays as  $e^{-\varepsilon d(\cdot, \operatorname{supp} f, V + |\mathbf{B}|)}$  away from the support of f and the resolvent decays as  $e^{-\varepsilon d(\cdot, \operatorname{supp} f, V + |\mathbf{B}| + \frac{1}{t^2})}$ .

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•  $L^2$  exponential decay for the resolvents has appeared in the literature, see Germinet and Klein [5], but purely in terms of  $\frac{1}{t^2}$ . Also in many other sources.

• The estimate (11) (for the operator  $L^{-1}$ ) is entirely new and is a consequence of the decay afforded by our assumptions on V and B.

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#### Definition 2

We say that the operator L has a *Moser estimate* if for each ball  $B \subset \mathbb{R}^n$  and each function u which solves Lu = 0 in the weak sense on 2B, it follows that  $u \in L^{\infty}(B)$  and

$$||u||_{L^{\infty}(cB)} \le C \Big( \int_{2B} |u|^2 \Big)^{\frac{1}{2}},$$
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where c, C are independent of B and u.

• The electric Schrödinger operators with real matrix *A*, and the magnetic Schrödinger operator both have Moser estimates.

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• The electric Schrödinger operators with real matrix *A*, and the magnetic Schrödinger operator both have Moser estimates.

#### Theorem 3 (Mayboroda, P. 2018)

Suppose  $\mathbf{a} \in L^2_{loc}(\mathbb{R}^n)$ , A is an elliptic matrix with complex, bounded coefficients,  $V \in L^1_{loc}(\mathbb{R}^n)$ , and that L is an operator for which there exists a fundamental solution bounded above by a multiple of  $|x-y|^{2-n}$ . Moreover, if  $\mathbf{a} \equiv 0$ , assume  $V \in RH_{\frac{n}{2}}$ ; otherwise assume (12). Then there exists  $\varepsilon > 0$  and a constant C > 0, depending on L, such that

$$\left(\int_{B(x,\frac{1}{m(x,V+|\mathbf{B}|)})} |\Gamma(z,y)|^2 \, dz\right)^{\frac{1}{2}} \le \frac{Ce^{-\varepsilon d(x,y,V+|\mathbf{B}|)}}{|x-y|^{n-2}} \text{ for all } x, y \in \mathbb{R}^n.$$

$$\tag{14}$$

If L has a Moser estimate, then

$$|\Gamma(x,y)| \le \frac{Ce^{-\varepsilon d(x,y,V+|\mathbf{B}|)}}{|x-y|^{n-2}} \text{ for all } x, y \in \mathbb{R}^n.$$
(15)

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Theorem 3 establishes in particular the upper bound exponential decay for both electric Schrödinger operators  $-\text{div }A\nabla + V$ , and the magnetic Schrödinger operator  $(\nabla - i\mathbf{a})^2 + V$ .

#### Definition 4

We say that the operator L satisfying assumptions (12) has the *m*-scale invariant Harnack Inequality if whenever  $B = B(x_0, r)$ ,  $r < \frac{c}{m(x_0, V + |\mathbf{B}|)}$ ,  $x_0 \in \mathbb{R}^n$ , the following property holds. For any u which solves Lu = 0 in the weak sense on 2B,

$$\sup_{x \in B} |u(x)| \le C \inf_{x \in B} |u(x)|, \tag{16}$$

with the constant C > 0 independent of B.

 $\bullet$  The electric Schrödinger operators with real matrix A have the  $m\mbox{-scale}$  invariant Harnack inequality.

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• The electric Schrödinger operators with real matrix A have the m-scale invariant Harnack inequality.

#### Theorem 5 (Mayboroda, P. 2018)

Suppose that  $\mathbf{a} \in L^2_{loc}(\mathbb{R}^n)$ , A is an elliptic matrix with complex, bounded coefficients,  $V \in L^1_{loc}(\mathbb{R}^n)$ , and that L,  $L_0 := L - V$ ,  $L_0^*$  are operators for which there exist fundamental solutions  $\Gamma \equiv \Gamma_V$ ,  $\Gamma_0 \Gamma_0^*$ . Assume that  $\Gamma_V, \Gamma_0, \Gamma_0^*$  are bounded above by a multiple of  $|x - y|^{2-n}$ , and that  $\Gamma_0$  is bounded below by a multiple of  $|x - y|^{2-n}$ . Suppose that L has a Moser estimate, and that L satisfies the m-scale invariant Harnack Inequality. Moreover, if  $\mathbf{a} \equiv 0$ , assume that  $V \in RH_{\frac{n}{2}}$ ; otherwise assume (12). Then there exist constants c and  $\varepsilon_2$  depending on  $\lambda, \Lambda, ||V + |\mathbf{B}||_{RH_{\frac{n}{2}}}$ , n and the constants from (12) such that

$$|\Gamma(x,y)| \ge \frac{ce^{-\varepsilon_2 d(x,y,V+|\mathbf{B}|)}}{|x-y|^{n-2}}.$$
(17)

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# Exponential decay of the real electric Schrödinger operators

If  $\Gamma_E$  is the fundamental solution to an electric Schrödinger operator with real matrix A and  $V \in RH_{\frac{n}{2}}$ , then the last few theorems imply that

$$\frac{c_1 e^{-\varepsilon_1 d(x,y,V)}}{|x-y|^{n-2}} \le \Gamma(x,y) \le \frac{c_2 e^{-\varepsilon_2 d(x,y,V)}}{|x-y|^{n-2}}.$$
(18)

Thank you. :)

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- S. Agmon, Lectures on exponential decay of solutions of second-order elliptic operators; bounds on eigenfunctions of N-body Schrödinger operators, Princeton Mathematical Notes 29 (1982).
- [2] D. Arnold, G. David, M. Filoche, D. Jerison, and S. Mayboroda, Localization of eigenfunctions via an effective potential, arXiv:1712.02419.
- B. Ben Ali, Maximal inequalities and Riesz transform estimates on L<sup>p</sup> spaces for magnetic Schrdinger operators I, Journal of Functional Analysis 259 (2010), 1631-1672.
- [4] C. Fefferman, *The uncertainty principle*, Bull. Amer. Math. Soc., 9 (1983), 129-206.
- [5] F. Germinet and A. Klein, Operator kernel estimates for functions of generalized Schrödinger operators, Proceedings of the American Mathematical Society 131 No. 3, (2002), 911-920.
- [6] K. Kurata, An estimate on the heat kernel of magnetic Schrödinger operators and uniformly elliptic operators with

*non-negative potentials*, J. London Math. Soc. **62** (2) (2000), 885-903.

- K. Kurata and S. Sugano, A remark on estimates for uniformly elliptic operators on weighted L<sup>p</sup> spaces and Morrey spaces, Math. Nachr. 209 (2000), 137–150.
- [8] V. Z. Meshkov, On the possible rate of decay at infinity of solutions of second order partial differential equations, Math. USSR-Sb., 72:2 (1992), 343361.
- Z. Shen, On fundamental solutions of generalized Schrödinger operators, Journal of Functional Analysis 167 (1999), 521-564.
- Z. Shen, Eigenvalue asymptotics and exponential decay of eigenfunctions for Schrödinger operators with magnetic fields, Trans. Amer. Math. Soc. 348 No.11 (1996), 4465-4488.
- Z. Shen, L<sup>p</sup> estimates for Schrödinger operators with certain potentials, Annales de l'institut Fourier 45 No.2 (1995), 513-546.
- Z. Shen, Estimates in L<sup>p</sup> for magnetic Schrödinger operators, Indiana University Mathematics Journal 45 No.3 (1996), 817-841.